

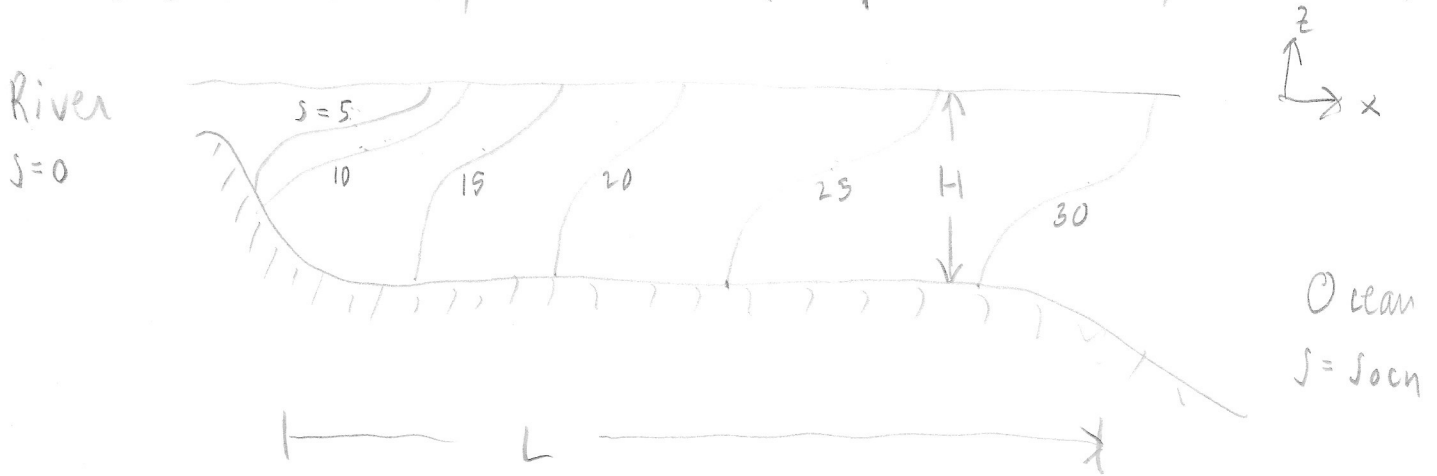
Estuarine Circulation

(7) 8/5/2019

(1)

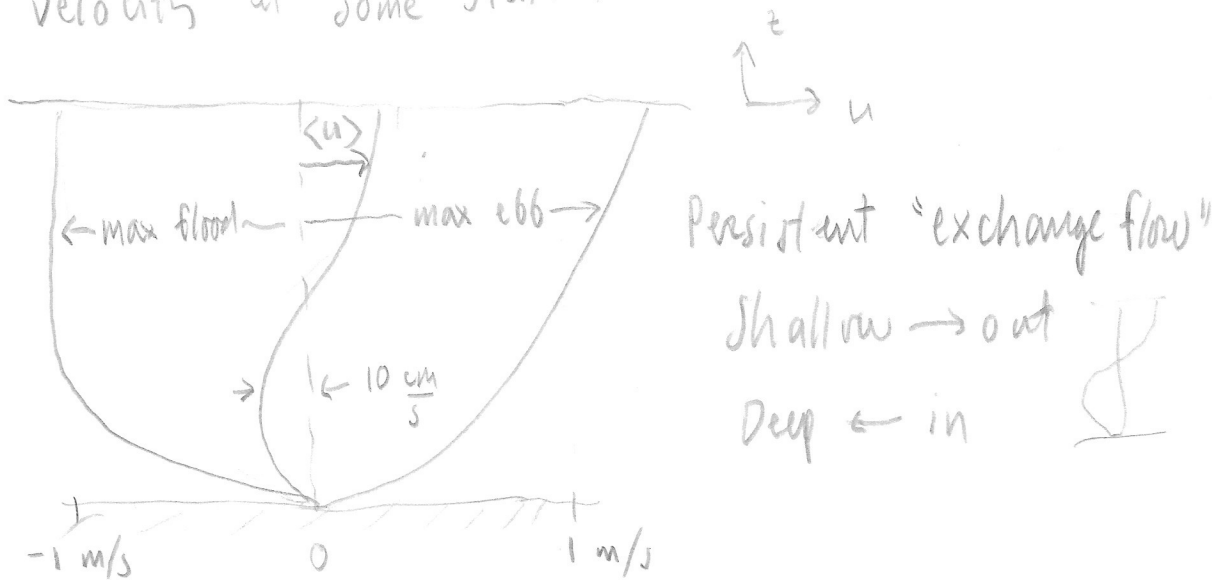
Estuary: a long bay influenced by rivers + tides

Observed salinity structure (examples in MacCreedy + Barnes 2011)

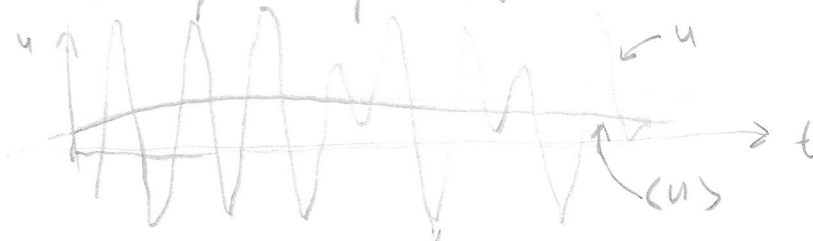


- Vertically stratified
- Along-channel density gradient

Observed velocity at some station:



$\langle u \rangle \equiv$ tidally averaged u ("subtidal", "residual")



Tidal-averaging: e.g. $\langle u \rangle_{t_0} =$ weighted average of 40-71 hours of u around t_0 (2)

e.g. LiveOcean / alpha / zfun_fit_godin()

• What causes exchange flow?

$$\langle \text{x mom} \rangle \quad \langle u \rangle_t + \langle \eta \cdot \nabla u \rangle - f \langle v \rangle = -\frac{1}{\rho_0} \langle p_x \rangle + \langle (A \cdot u_z)_z \rangle$$

\swarrow steady exchange
 \swarrow wave advection is small
 \swarrow narrow channel $\langle v \rangle \approx 0$
 $\underbrace{\hspace{10em}}$ pressure gradient balanced by vertical stress divergence

pressure: hydrostatic $\Rightarrow p = g \int_z^\eta \rho dz$ and $p = p_0 (1 + \beta s)$

$$\Rightarrow p_x \approx \rho_0 g \eta_x + g \int_z^0 \rho_0 \beta s_x dz$$

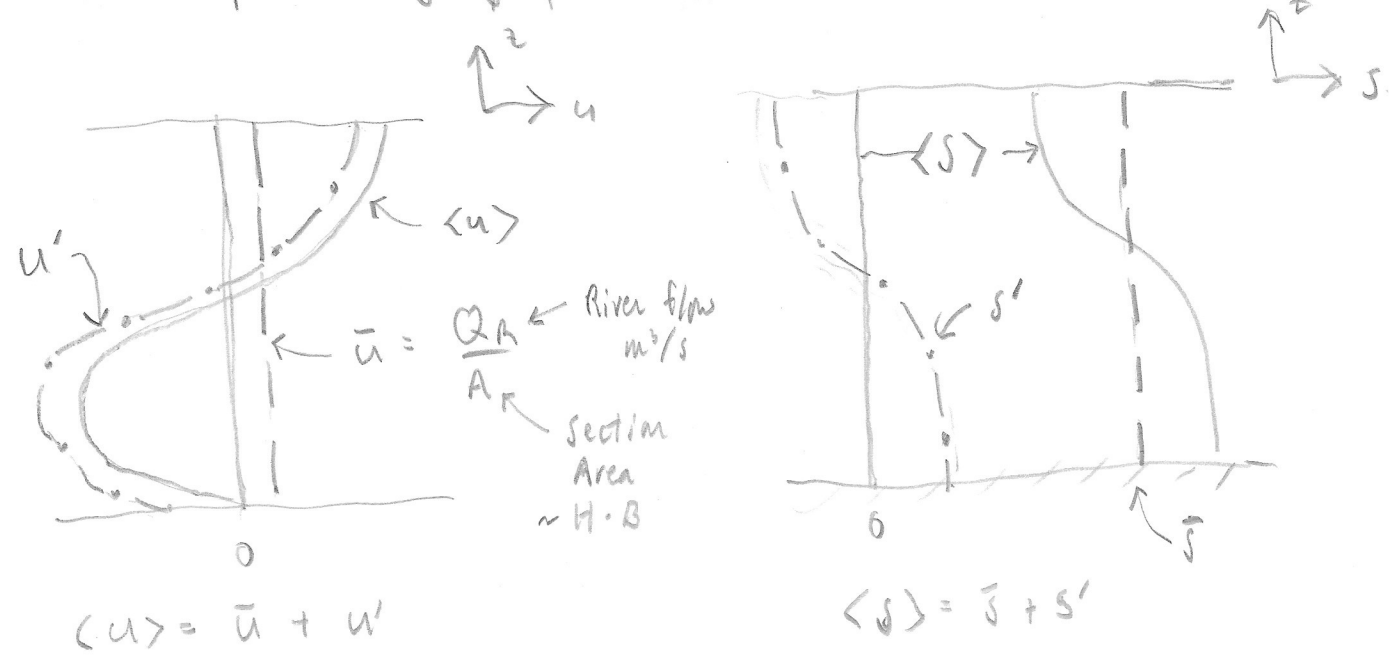
$$\beta = 7.7 \times 10^{-4}$$

$$\rho_0 = 1000 \text{ kg/m}^3$$

and $-\frac{\langle p_x \rangle}{\rho_0} = -g \langle \eta_x \rangle - g \int_z^0 \beta \langle s_x \rangle dz$

We are only concerned with tidally-averaged flow so drop the $\langle \rangle$ when writing variables.

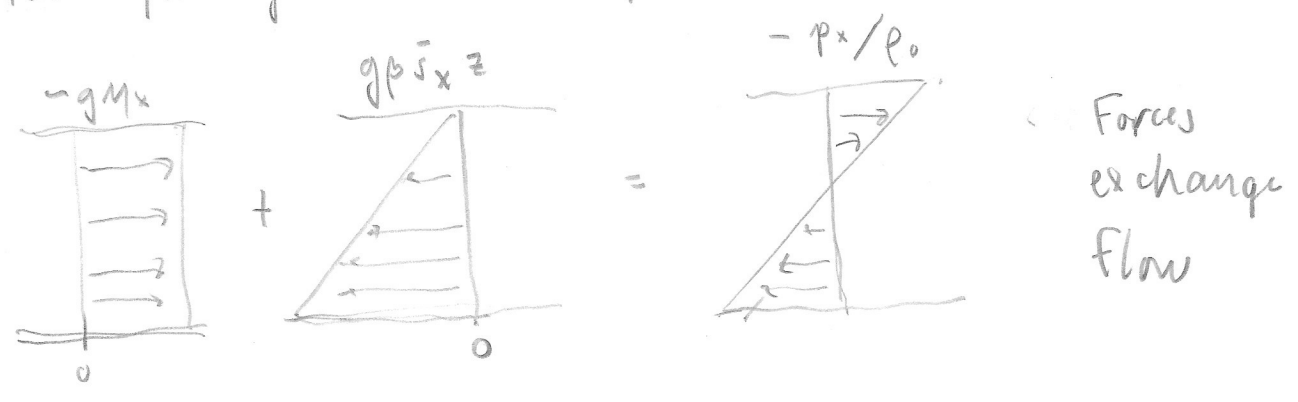
And we decompose into vertically-averaged and depth-varying parts:



We observe (Pritchard 1954): $[s'_x] \ll [\bar{s}_x]$

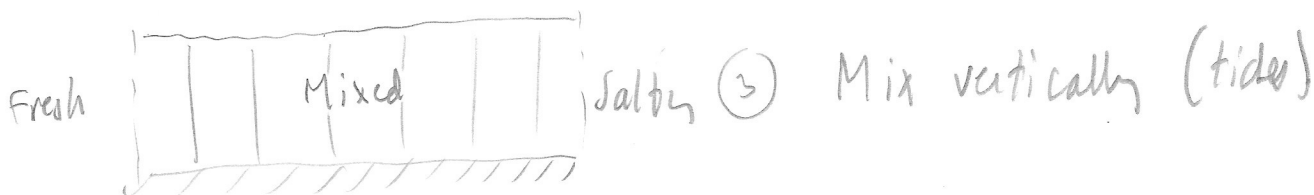
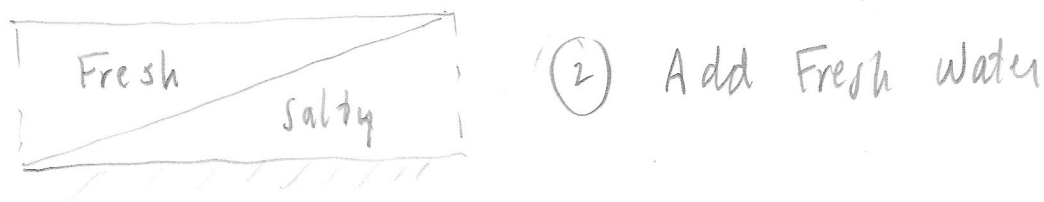
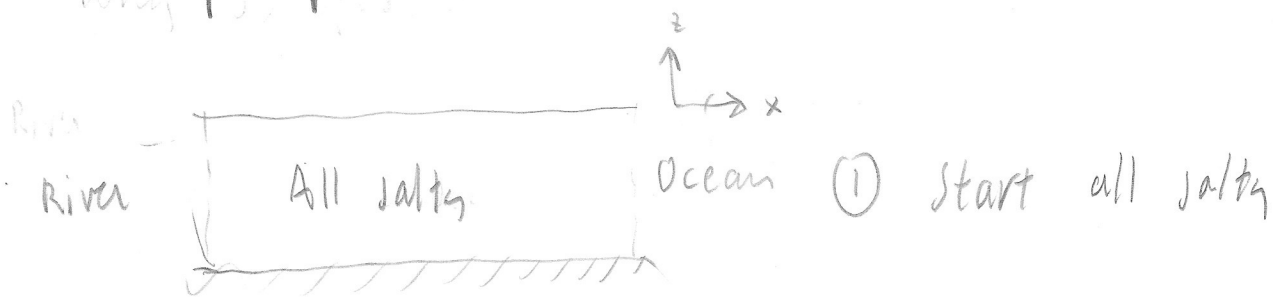
so $-\frac{p_x}{\rho_0} = -g \eta_x + g \beta \bar{s}_x z$ (dropped $\langle \rangle'$)

For η_x negative and \bar{s}_x positive:



Thought problem about \bar{S}_x :

(4)



$\Rightarrow \bar{S}_x > 0$

- Then
- Shallow exchange flow removes mixed water
 - river adds FW
 - deep exchange flow adds salt water (back to (1)/(2))

In reality all are happening at once, or system is going (2) \leftrightarrow (3) over spring-neap cycle.

We are doing math to find $w(z)$ given \bar{S}_x

To find $u(z)$ we solve

assume $A = \text{const.}$

(5)

$$0 = -g\gamma_x + g\beta \bar{s}_x z + A u_{zz}$$

subject to boundary conditions

(i) $u(-H) = 0$

No slip at bottom

(ii) $u_z(0) = 0$

No stress at top

and integral constraint:

(iii) $\frac{1}{H} \int_{-H}^0 u dz = \bar{u}$

note: $\int_{-H}^0 u' dz = 0$

A , \bar{s}_x and \bar{u} are assumed known

so we integrate twice in z , using (i) + (ii)

and use (iii) to find right value of γ_x .

Result is cubic in z :

Result

$$u = \bar{u} + \left[\bar{u} \frac{1}{z} (1 - 3\eta^2) + u_E (1 - 9\eta^2 - 8\eta^3) \right]$$

where $\eta = z/H$ and $u_E = \frac{g\beta \bar{s}_x H}{48} \frac{H^2}{A}$

u_E expresses balance:

pressure gradient

vertical friction time scale

U_E is the scale of the exchange flow

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Note $U_E \sim H^3$

\Rightarrow exchange much stronger for deep estuaries?

No: because \bar{S}_x adjusts dynamically depending on how much mixing happens

So: we need to consider controls $m(\delta)$.

(*) For the frictionless case, if we allow acceleration: $U_t = -g\eta_x + g\beta\bar{S}_x z$

• If $\bar{S}_x = 30/50 \text{ km}$, $H = 20 \text{ m}$, and $\eta_x = 0$

how fast is the bottom water going after 3 hours?

= 10 m/s

• what value of η_x would be required to oppose the acceleration due to \bar{S}_x so that $\int_{-H}^0 U_t dz = 0$?

My answers:

$$u_t = -g \eta_x + g \beta \bar{s}_x z$$

assume 0

$\frac{30}{10^5} \frac{1}{m} = \frac{\Delta s}{L}$

-H

$$\Rightarrow u = -g \beta \bar{s}_x H t = -\frac{1}{10} \cdot 7.7 \times 10^{-4} \cdot \frac{50}{10^5} \cdot 20 \cdot 3600 \cdot 3$$

$$= 7.7 \times 10^{-4} \cdot \frac{50 \cdot 2 \cdot 3.6 \cdot 3}{10^5} = 0.8 \text{ m/s}$$

= u(-H, 3 hours)

$$\int_{-H}^0 \dot{u}_t dz = 0 = -g \eta_x H + g \beta \bar{s}_x \left(\frac{1}{2} z^2 \right) \Big|_{-H}^0$$

$$g \eta_x H = -\frac{1}{2} g \beta \bar{s}_x H^2$$

$$\frac{\Delta \eta}{x} = -\frac{1}{2} \beta \frac{\Delta s}{x} H = -\frac{1}{2} 7.7 \times 10^{-4} \times 50 \times 20$$

$$= -\frac{1}{2} 7.7 \times 10^{-4} \approx -0.4 \text{ m} !$$

= Δη


But observations suggest η_x is much smaller than this.

Geyer et al. (2000 JPO) Hudson River Estuary

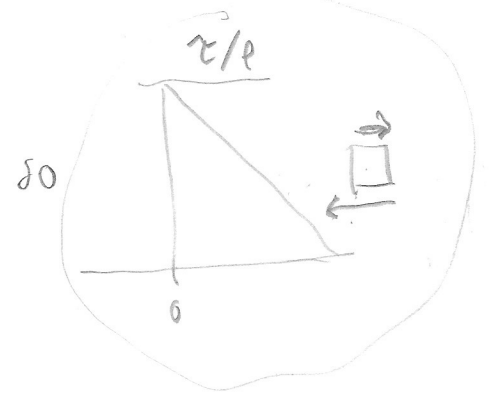
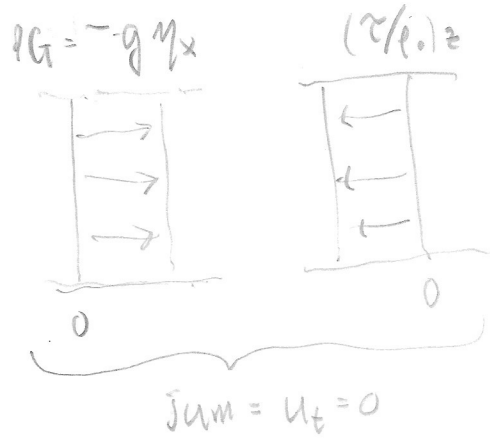
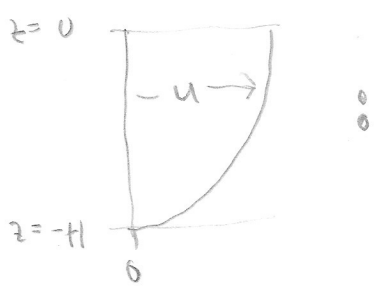
Looking at the force balance :

x mm $U_t = \underbrace{-\frac{1}{\rho_0} \rho_x}_{PG} + \frac{\partial}{\partial z} (\tau/\rho_0)$, $\frac{\text{stress}}{\rho_0} = \frac{\tau}{\rho_0} = A \frac{\partial u}{\partial z} = -\langle u'w' \rangle$

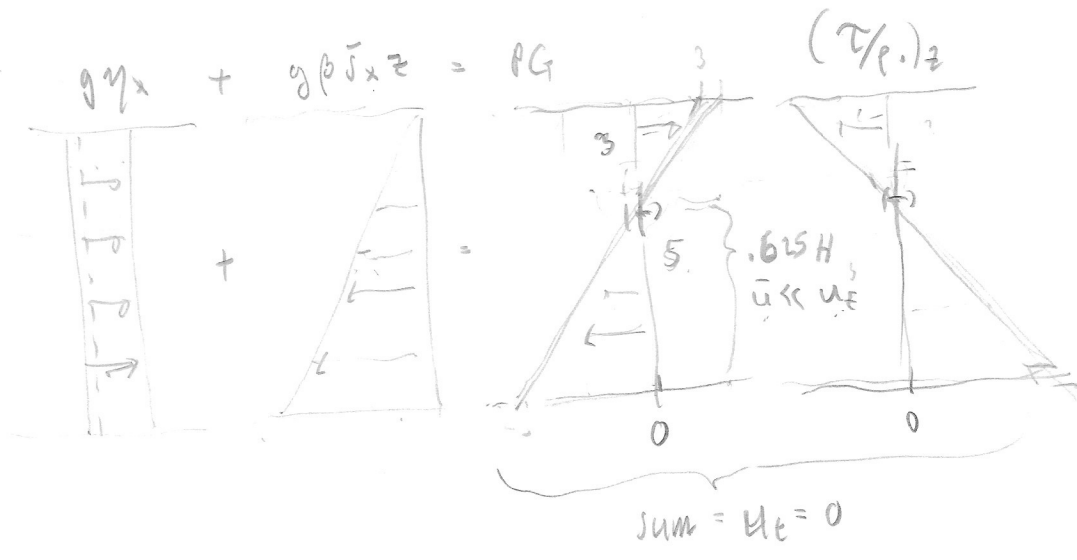
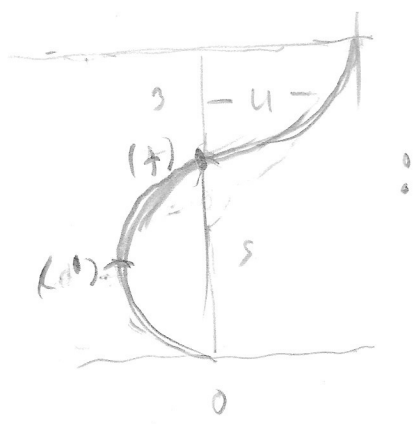
steady

sign convention 

unidirectional:

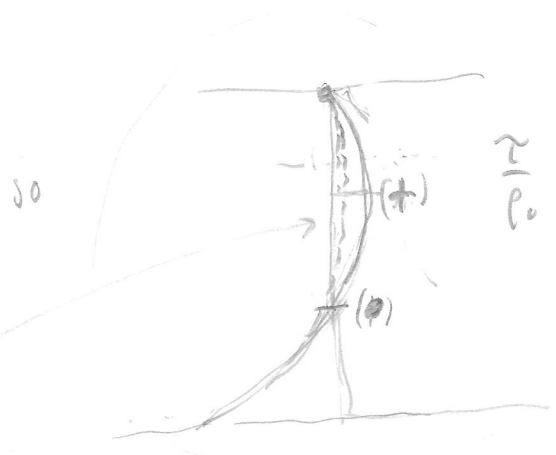


exchange flow



In observations in Hudson River
(Geyer et al. 2000 JLO)

the stress in the surface
layer, and η_x , are
much smaller



Appendix What is the Au_{zz} profile for our simple cube u ? (9)

8/6/2019

$$u_t = -g\eta_x + g\beta\bar{s}_k z + Au_{zz} \quad u_E$$

$$u = \bar{u} + \bar{u} \frac{1}{2}(1-3\eta^2) + u_E (1-9\eta^2 - 8\eta^3)$$

$$\eta = z/H$$

$$2\eta = \frac{1}{H} z$$

$$d\eta = \frac{1}{H} dz$$

$$u_{zz} = \bar{u}(-3)2\frac{z}{H} + u_E \left[(-9)(2)\frac{z}{H} - 8(3)\frac{z^2}{H^2} \right]$$

$$\frac{\partial}{\partial z} = \frac{1}{H} \frac{\partial}{\partial \eta}$$

$$u_{zz} = -6\bar{u}\frac{z}{H} + u_E \left(-18\frac{z}{H} - 24\frac{z^2}{H^2} \right)$$

$$u_{zz} = -\frac{6\bar{u}}{H^2} + u_E \left(-\frac{18}{H^2} - \frac{48z}{H^2} \right)$$

$$u_{zz}|_0 = -\frac{6\bar{u}}{H^2} + u_E \left(-\frac{18}{H^2} \right)$$

$$u_{zz}|_{-H} = -\frac{6\bar{u}}{H^2} + u_E \left(-\frac{18}{H^2} + \frac{48}{H^2} \right) = -\frac{6\bar{u}}{H^2} + u_E \frac{30}{H^2}$$

$\eta = -1 \rightarrow$

$$Au_{zz}|_0 = \frac{A}{H^2} (-6\bar{u} - 18u_E)$$

$$Au_{zz}|_{-H} = \frac{A}{H^2} (-6\bar{u} + 30u_E)$$

